

**A LOOK AT THE DUFFING SYSTEM WITH PERIODIC  
PARAMETRIC EXCITATION: AN OVERVIEW  
(WEB-PLATFORM UPGRADE)**

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**ABSTRACT:** In this article we demonstrate some specialized modules for investigating the dynamics of some generalized Duffing system with periodic parametric excitation, an integral part of a planned much more general Web-based application for scientific computing. We also study some new hypothetical oscillators. Numerical examples, illustrating our results using *CAS MATHEMATICA* are given.

**Key Words:** generalized Duffing system with periodic parametric excitation, Duffing oscillator, hypothetical oscillator, diagram factor

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## 1. INTRODUCTION

In this paper we demonstrate some specialized modules for investigating the dynamics of some generalized Duffing system with periodic parametric excitation, an integral

part of a planned much more general Web-based application for scientific computing (for some details see [10]–[17]).

We also study some new hypothetical oscillators.

More precisely, this WEB Platform envisages research on: Duffing system with periodic parametric excitation; modified Duffing systems; new hypothetical oscillators; a modification of the basic Duffing system with periodic parametric excitation.

Numerical examples, illustrating our results using *CAS MATHEMATICA* are also given.

In [7]–[8] some investigations and simulations on the planar Rayleigh–Lienard system are given.

We note that the use of normalized diagram factor  $\frac{y(b \cos \theta + c)}{N}$  (where  $\theta$  is the azimuthal angle and  $c$  is the phase difference) is very complicated.

Where possible, the corresponding diagram-functions (with application in the field of antenna analysis and synthesis) generated by the model oscillators investigated in this article have been generated and visualized.

A natural extension of the planned much more general Web-based application for scientific computing involves further consideration and simulations on generalized Duffing systems with periodic parametric excitation.

## 2. MAIN RESULTS. SIMULATIONS

### 2.1. A LOOK AT THE DUFFING SYSTEM WITH PERIODIC PARAMETRIC EXCITATION

The Rayleigh–Duffing oscillator models are widely used in physics, electronics, and many other disciplines.

Consider the following planar system (see for example [1])

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = f \cos(gt)x - x^3 - hy \end{cases} \quad (1)$$

where  $h$  is the damping coefficient, and  $f$  and  $g$  are the amplitude and frequency of excitation.

**I.** For given  $f = 6.25$ ,  $g = 0.1$ ;  $h = 0.3$ , the simulations on the system (1) for  $x_0 = 0.2$ ;  $y_0 = 0.1$  are depicted on Fig. 1.

**I.1** For given  $f = 5$ ,  $g = 0.05$ ;  $h = 0.01$ , the simulations on the system (1) for  $x_0 = 0.3$ ;  $y_0 = 0.2$  are depicted on Fig. 2.

## 2.2. A LOOK AT THE MODIFIED DUFFING SYSTEM

Chaotic motions of a Rayleigh–Duffing oscillator with periodically external and parametric excitations are investigated rigorously.

For some results see [24]–[25].

Consider the planar system (see for example [2])

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = x - x^3 + \epsilon(a_1 \cos(at) - a_2 y) \end{cases} \quad (2)$$

where  $0 \leq \epsilon \leq 1$ .

**II.** For given  $a_1 = 6$ ,  $a_2 = 2$ ,  $a = 0.005$ ,  $\epsilon = 0.005$ , the simulations on the system (2) for  $x_0 = 0.2$ ;  $y_0 = 0.1$  are depicted on Fig. 3.

**II.1** For given  $a_1 = 0.1$ ,  $a_2 = 0.003$ ,  $a = 0.005$ ,  $\epsilon = 0.1$ , the simulations on the system (2) for  $x_0 = 0.3$ ;  $y_0 = 0.2$  are depicted on Fig. 4.

For fixed  $b = 0.92$ ,  $c = 0.73$  the normalized diagram factor is depicted on Fig. 4 d).

Remark. In [6] the authors study the dynamics of a generalized oscillator model on the base of model (2) in the light of Melnikov’s approach.

Consider the planar system:

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = a_1 x - a_2 x^3 + \epsilon((1 - y^2)y + (1 + x) \cos(gt)) \end{cases} \quad (3)$$

where  $0 \leq \epsilon \leq 1$ .

**III.** For given  $a_1 = 0.2$ ,  $a_2 = 0.9$ ,  $g = 0.02$ ,  $\epsilon = 0.01$ , the simulations on the system (3) for  $x_0 = 0.3$ ;  $y_0 = 0.2$  are depicted on Fig. 5.

For fixed  $b = 0.9$ ,  $c = 0.45$  the normalized diagram factor is depicted on Fig. 5 d).

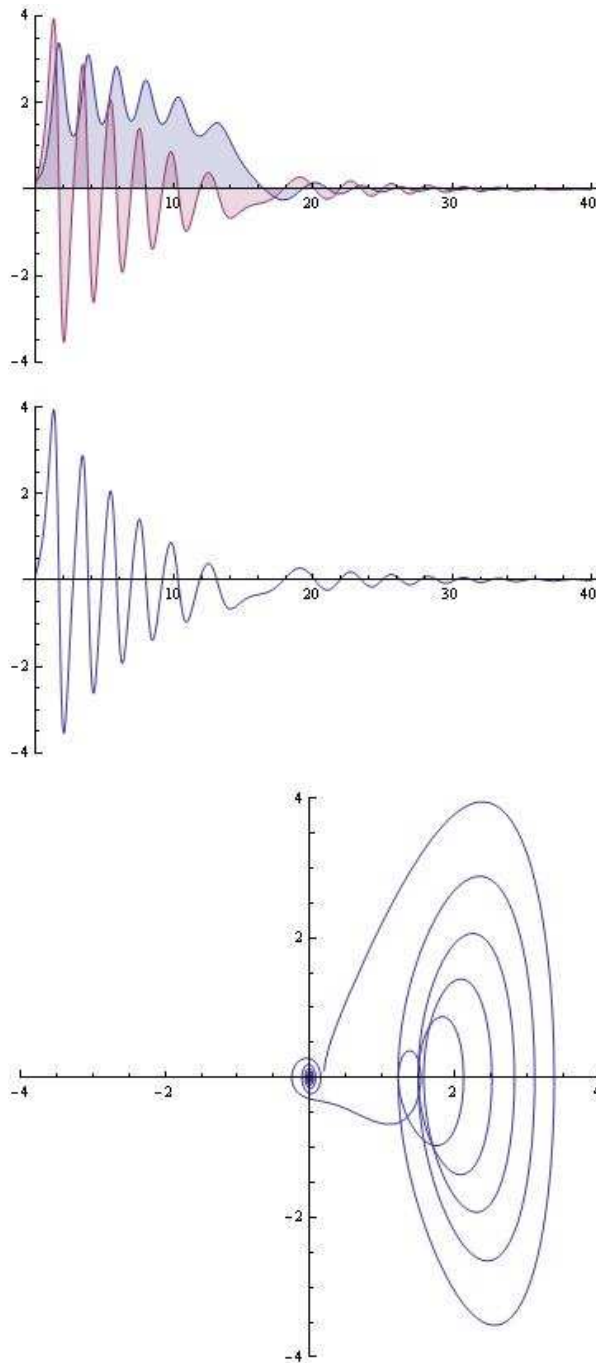


Figure 1: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait (example I).

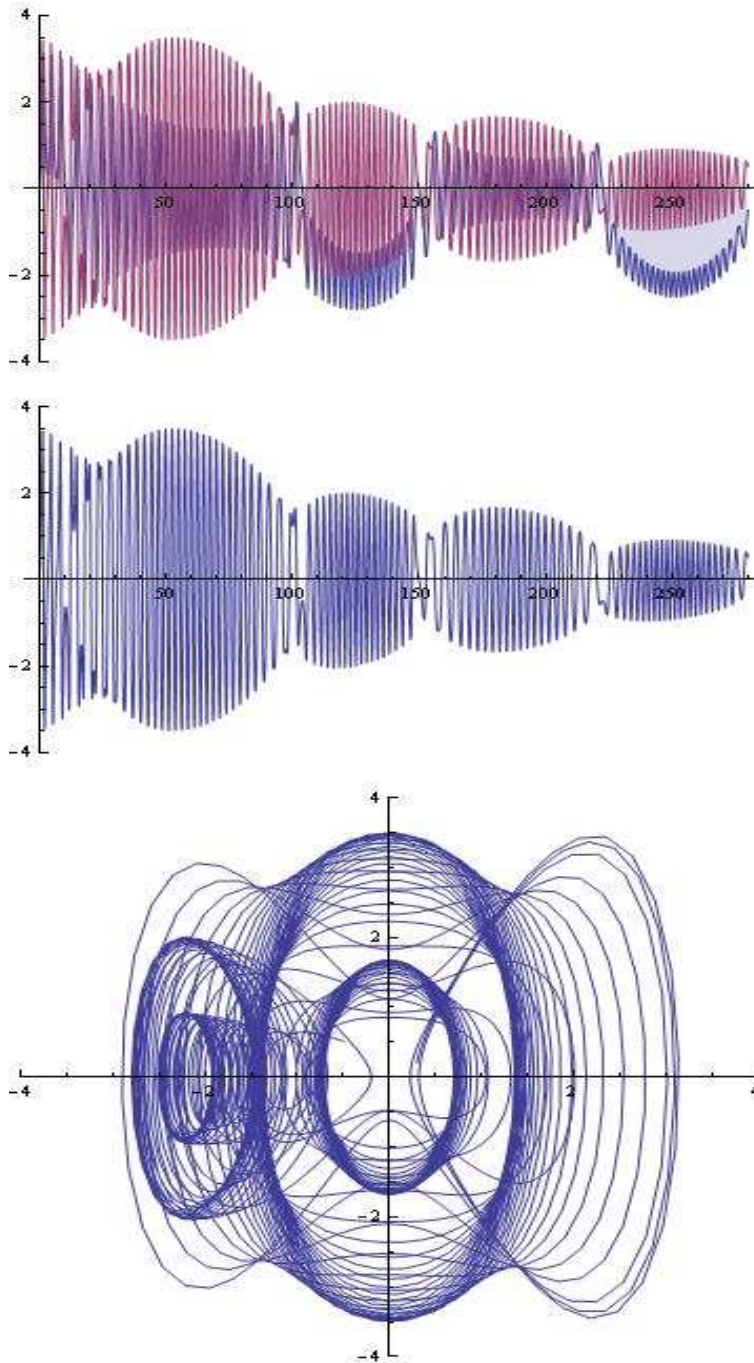


Figure 2: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait (example I.1).

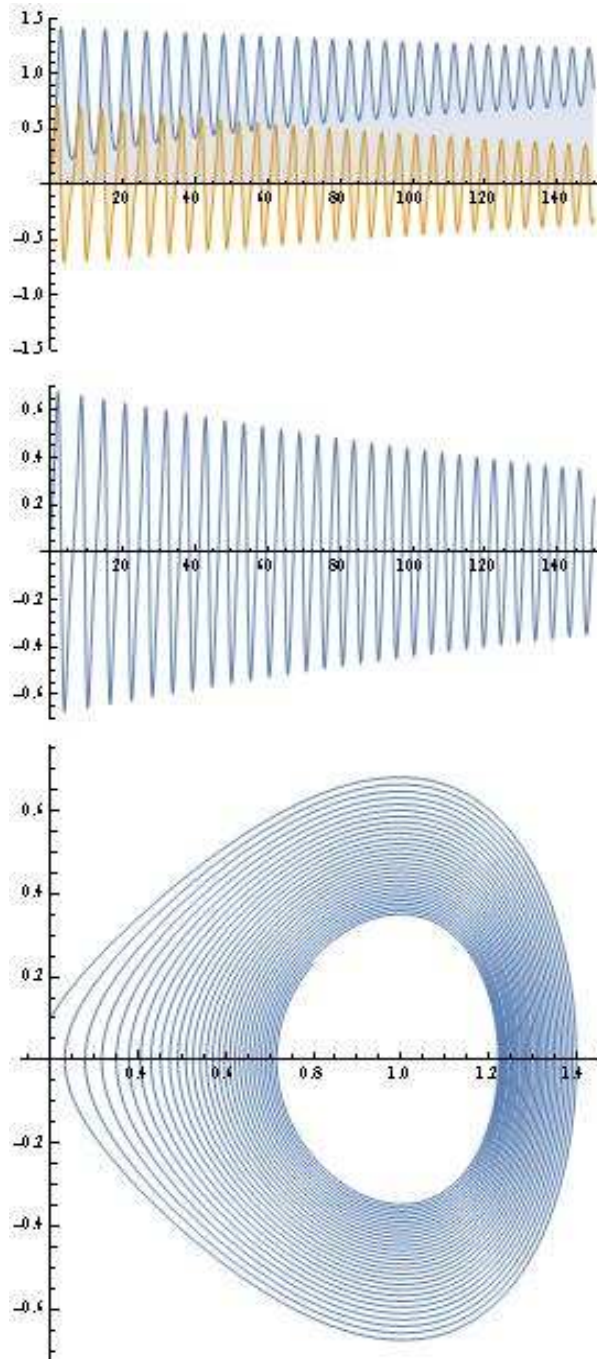


Figure 3: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait (example II).

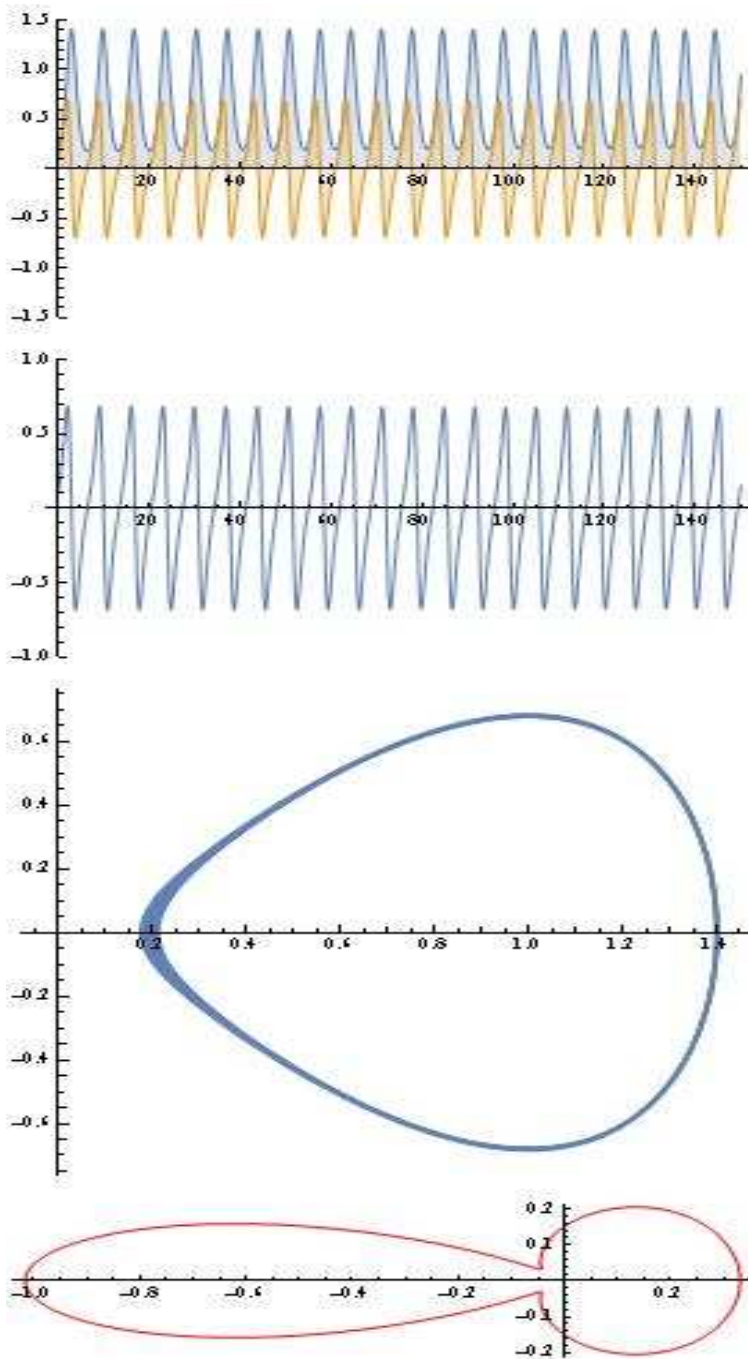


Figure 4: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait; d) diagram factor (example II.1).

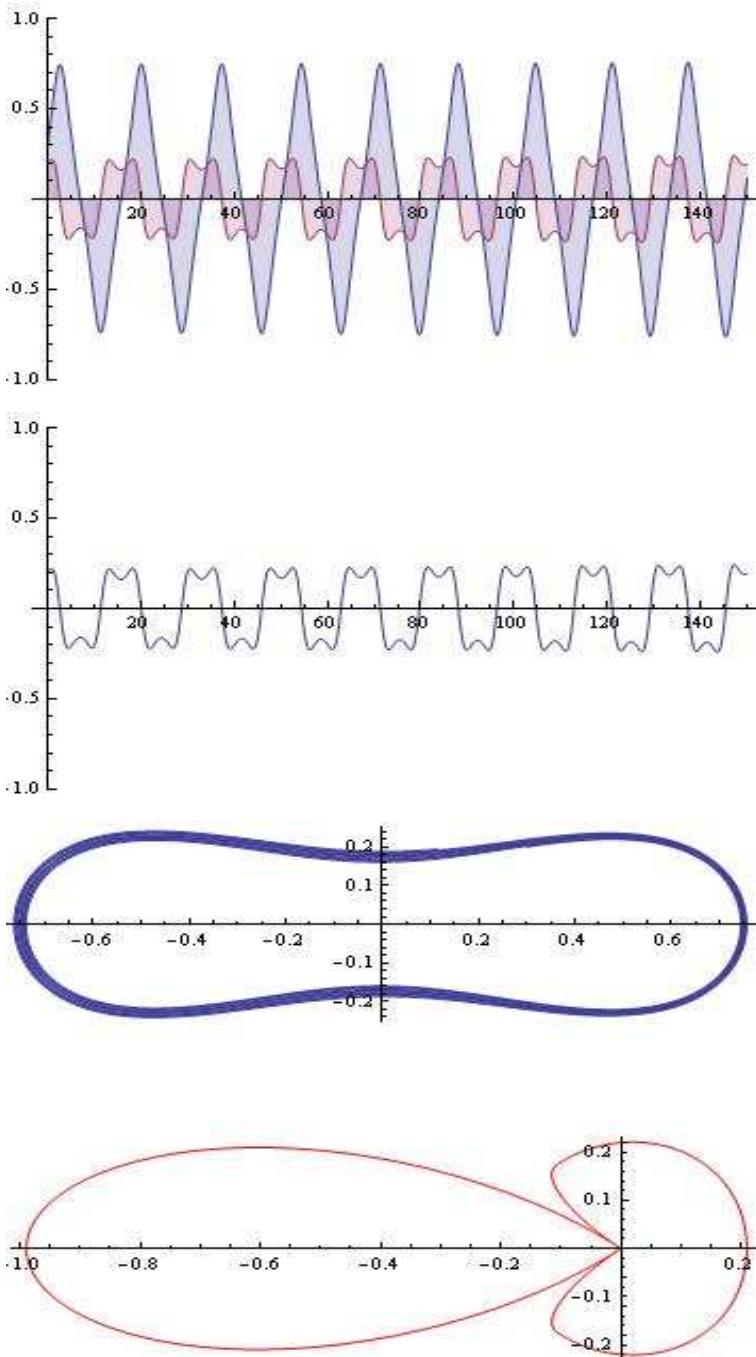


Figure 5: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait; d) diagram factor (example III).



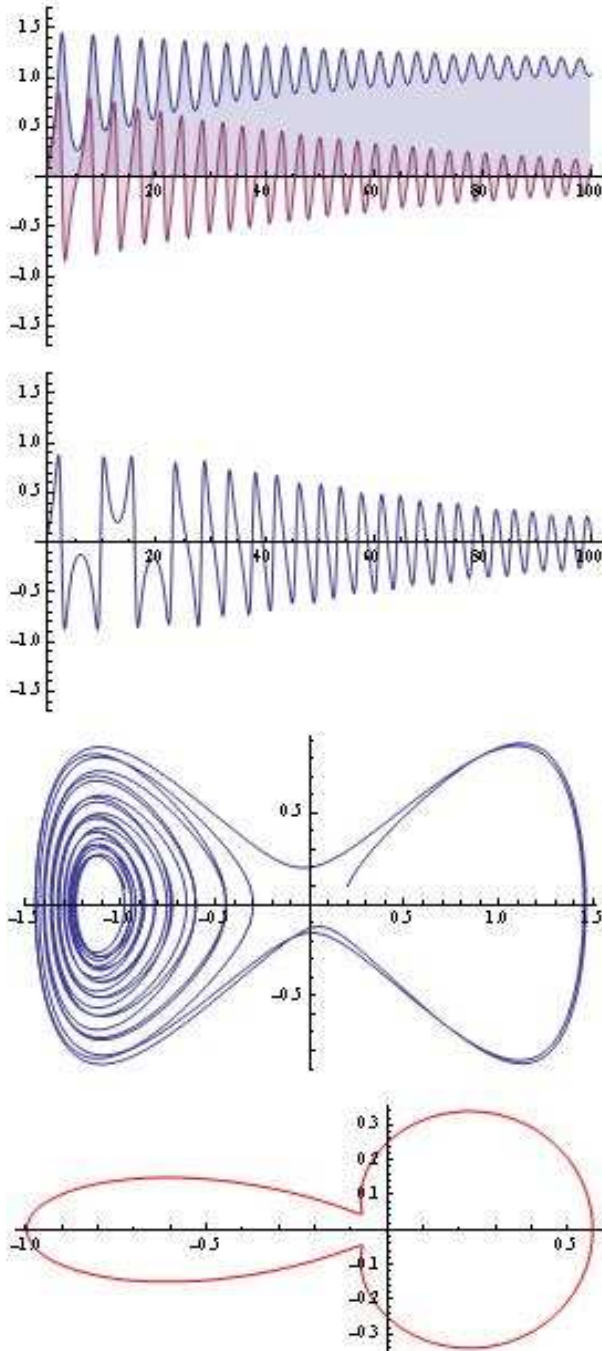


Figure 6: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait; d) diagram factor (example IV).

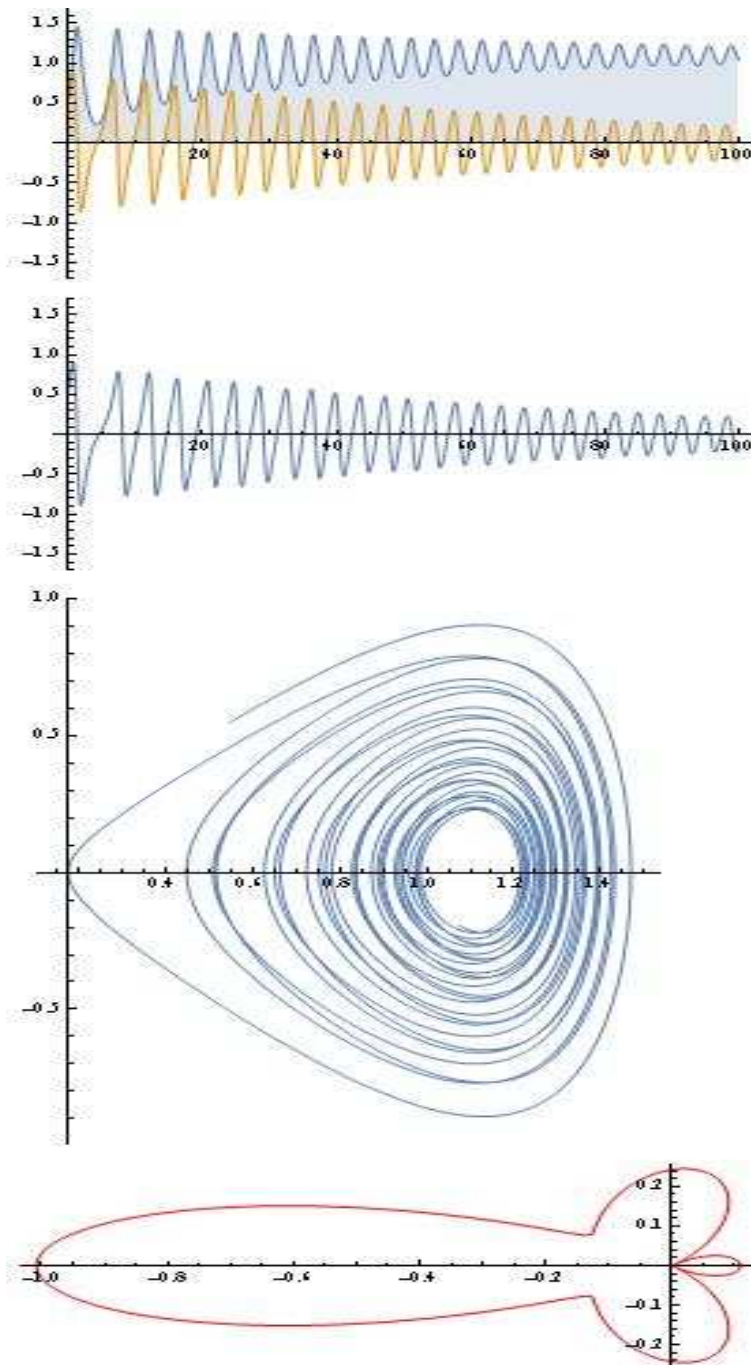


Figure 7: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait; d) diagram factor (example IV.1).

### 2.3. A LOOK AT THE NEW HYPOTHETICAL OSCILLATORS

Consider the following planar system

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = x - \sum_{i=0}^{\lfloor \frac{n}{2} \rfloor - 1} \frac{x^{n-2i}}{n-2i} + \epsilon(a_1 \cos(at) - a_2 y) \end{cases} \quad (4)$$

where  $0 \leq \epsilon \leq 1$ .

Remark. Our considerations on the Melnikov homoclinic integral are summarized in Appendix 1.

**IV.** For given  $n = 7$ ,  $a_1 = 0.25$ ,  $a_2 = 0.35$ ,  $a = 0.55$ ,  $b = 0.81$ ,  $c = 0.66$ ,  $\epsilon = 0.005$ , the simulations on the system (4) for  $x_0 = 0.2$ ;  $y_0 = 0.1$  are depicted on Fig. 6.

**IV.1** For given  $n = 7$ ,  $a_1 = 0.2$ ,  $a_2 = 0.3$ ,  $a = 0.5$ ,  $b = 0.8$ ,  $c = 0.665$ ,  $\epsilon = 0.005$ , the simulations on the system (4) for  $x_0 = 0.55$ ;  $y_0 = 0.55$  are depicted on Fig. 7.

### 2.4. A MODIFICATION OF THE MODEL (1)

Consider the following modification of model (1):

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = f \cos(gt)x - \sum_{i=0}^{\lfloor \frac{n}{2} \rfloor - 1} x^{n-2i} - hy \end{cases} \quad (5)$$

**V.** For given  $n = 7$ ,  $f = 4$ ,  $g = 0.07$ ,  $h = 0.03$ , the simulations on the system (5) for  $x_0 = 0.4$ ;  $y_0 = 0.3$  are depicted on Fig. 8.

**V.1** For given  $n = 9$ ,  $f = 5$ ,  $g = 0.005$ ,  $h = 0.01$ , the simulations on the system (5) for  $x_0 = 0.3$ ;  $y_0 = 0.2$  are depicted on Fig. 9.

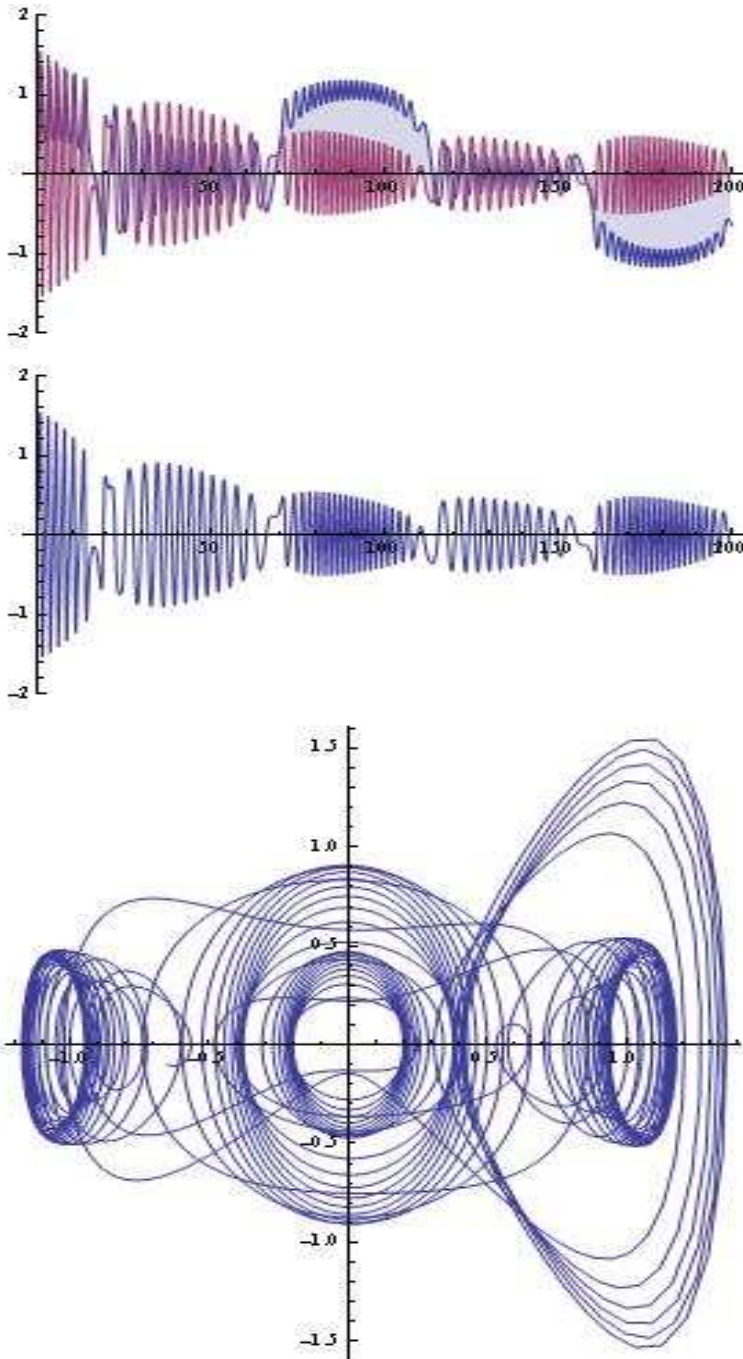


Figure 8: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait (example V).

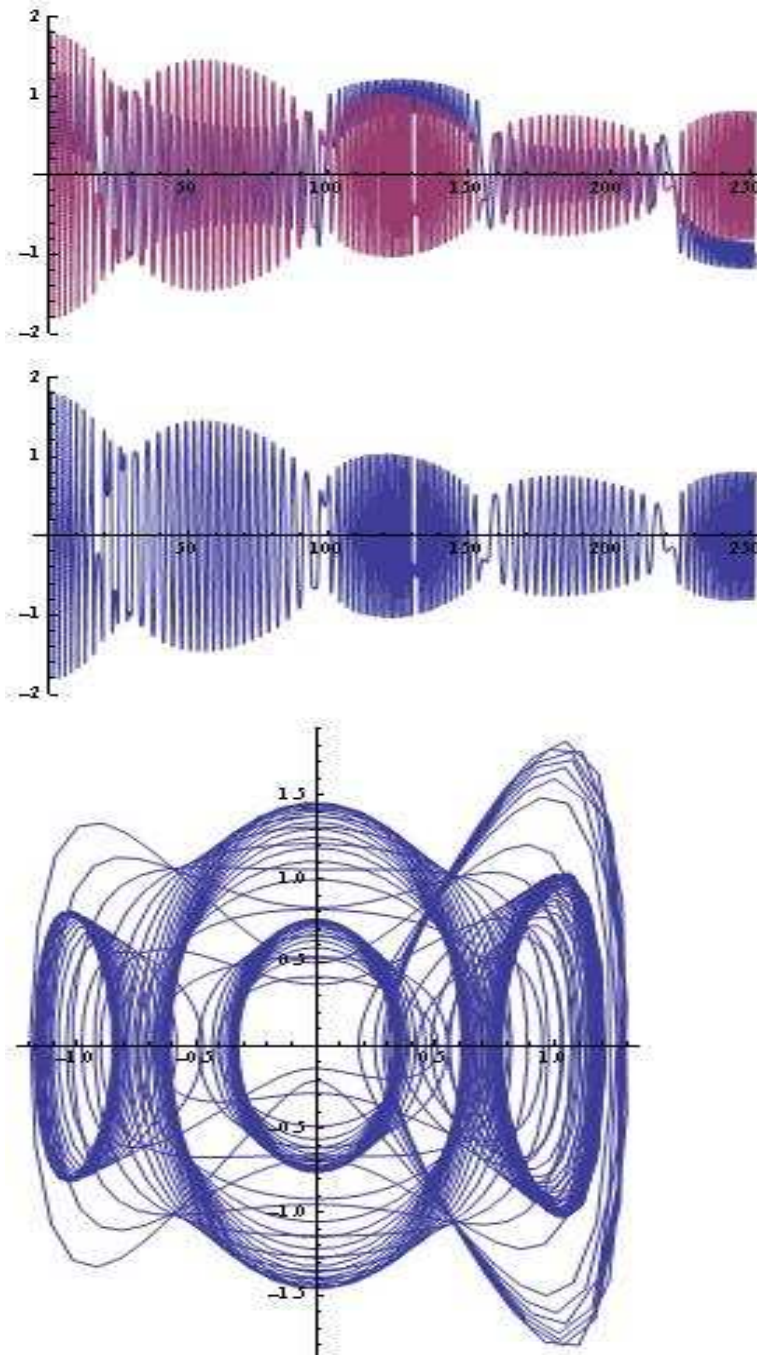


Figure 9: a) The solutions of differential system; b)  $y$ -component of the solution; c) Phase portrait (example V.1).

### 3. CONCLUDING REMARKS

In this article we demonstrate some specialized modules for investigating the dynamics of some generalized Duffing system with periodic parametric excitation, an integral part of a planned much more general Web-based application for scientific computing. More precisely, this WEB Platform envisages research on: Duffing system with periodic parametric excitation; modified Duffing systems; new hypothetical oscillators; a modification of the basic Duffing system with periodic parametric excitation and others. Where possible, we employ various optimization techniques for highperformance calculations, including multi-processor and multi-threading calculations, and hardware intrinsics [15]–[17]. We presented only a small part of the platform's capabilities. We will be grateful to all colleagues who, with their critical remarks, will contribute to its significant improvement. We fully understand that the construction of such an ambitious Web-based platform for scientific computing can only be realized with the active participation of specialists from various branches of scientific knowledge.

Remark. The study of corresponding critical levels of  $H(x, y) = \frac{1}{2}y^2 - P(x)$  is very complicated.

In this regard, we recommend the excellent study by Gavrilov and Iliev [28].

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### 4. APPENDIX 1.

The case  $n = 3$ . The system of the type (4)

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = x - \frac{1}{3}x^3 + \epsilon(a_1 \cos(at) - a_2y) \end{cases}$$

has the following Hamiltonian ( $\epsilon = 0$ )

$$H(x, y) = \frac{1}{2}y^2 - \frac{1}{2}x^2 + \frac{1}{12}x^4.$$

The homoclinic orbit is given by (see Fig.10)

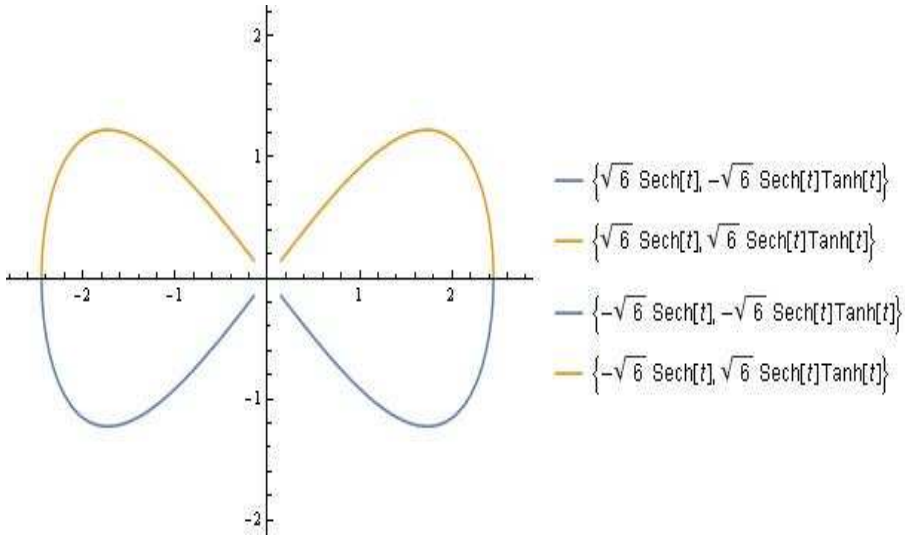


Figure 10: The homoclinic orbit.

$$x_0(t) = \pm\sqrt{6}sech t$$

$$y_0(t) = \mp\sqrt{6}sech t tanh t.$$

The Melnikov homoclinic integral is given by:

$$M(t_0) = \int_{-\infty}^{\infty} y_0(t) (a_1 \cos(a(t + t_0)) - a_2 y_0(t)) dt.$$

From a numerical point of view, the task of finding a multiple root of  $M(t_0)$  is more interesting given that the parameters appearing in the proposed differential model are subject to a number of restrictions of a physical nature.

The following is valid

**Proposition 1.** For  $n = 3$  and  $D = K$  where

$$D = \frac{4a_2}{\sqrt{6}\pi} cosh \left( \frac{a\pi}{2} \right); K = aa_1$$

the Melnikov function  $M(t_0)$  has root with multiplicity two.

Proof. We have

$$\begin{aligned}
M(t_0) &= \int_{-\infty}^{\infty} y_0(t) (a_1 \cos(a(t + t_0)) - a_2 y_0(t)) dt \\
&= \int_{-\infty}^{\infty} (-a_2 y_0(t)^2 + y_0(t) a_1 \cos(a(t + t_0))) dt \\
&= -4a_2 + \sqrt{6} a a_1 \pi \operatorname{sech} \left( \frac{a\pi}{2} \right) \sin(at_0) \\
&= \sqrt{6} a a_1 \pi \operatorname{sech} \left( \frac{a\pi}{2} \right) \left( \sin(at_0) - \frac{D}{K} \right)
\end{aligned}$$

This completes the proof of Proposition 1.

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