

BLOW UP FOR COUPLED NONLINEAR WAVE EQUATIONS WITH WEAK DAMPING TERMS

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Abstract: In this paper, we consider a coupled nonlinear wave equations with weak damping terms. We prove the blow up of solutions with negative initial energy.

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1. Introduction

In this paper, we consider the following initial-boundary value problem

$$\left\{ \begin{array}{ll} u_{tt} + |u_t|^{p-1} u_t = \operatorname{div} \left(\rho \left(|\nabla u|^2 \right) \nabla u \right) + f_1(u, v), & (x, t) \in \Omega \times (0, T), \\ v_{tt} + |v_t|^{q-1} v_t = \operatorname{div} \left(\rho \left(|\nabla v|^2 \right) \nabla v \right) + f_2(u, v), & (x, t) \in \Omega \times (0, T), \\ u(x, t) = v(x, t) = 0, & (x, t) \in \partial\Omega \times (0, T), \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), & x \in \Omega, \\ v(x, 0) = v_0(x), \quad v_t(x, 0) = v_1(x), & x \in \Omega \end{array} \right. \quad (1.1)$$

where Ω is a bounded domain with smooth boundary $\partial\Omega$ in R^n , $n = 1, 2, 3$;

$p, q \geq 1$; $f_i : R^2 \rightarrow R$ are given functions to be specified later. Problems of this type arise in material science and physics.

Throughout this paper, we define ρ by

$$\rho(s) = b_1 + b_2 s^m, \quad q \geq 0, \quad (1.2)$$

where b_1, b_2 are nonnegative constants, and $b_1 + b_2 > 0$. Set $P(s) = \int_0^s \rho(\xi) d\xi$, $s \geq 0$.

(A1). Let $F(u, v) = a|u + v|^{r+1} + 2b|uv|^{\frac{r+1}{2}}$ with $a, b > 0$, $r \geq 3$ if $n = 1, 2$ and $r = 3$ if $n = 3$; $f_1(u, v) = \frac{\partial F}{\partial u}$, $f_2(u, v) = \frac{\partial F}{\partial v}$; $p, q \geq 1$ if $n = 1, 2$ and $1 \leq p, q \leq 5$ if $n = 3$.

One can easily verify that

$$u f_1(u, v) + v f_2(u, v) = (r + 1) F(u, v), \quad \forall (u, v) \in R^2. \quad (1.3)$$

Lemma 1. (see [2]) *There exist two positive constants c_0 and c_1 such that*

$$c_0 \left(|u|^{r+1} + |v|^{r+1} \right) \leq F(u, v) \leq c_1 \left(|u|^{r+1} + |v|^{r+1} \right)$$

is satisfied.

Wu et al. [6] obtained the global existence and blow up of the solution of problem (1.1) under some suitable conditions. Fei and Hongjun [1] considered problem (1.1) and improved the blow up result obtained in [6], for a large class of initial data in positive initial energy, using the some techniques as in Payne and Sattinger [4] and some estimates used firstly by Vitillaro [5]. Recently, Pişkin and Polat [3] studied the local and global existence, energy decay and blow of the solution of problem (1.1). Also, for more information about (1.1), see references [1, 3].

In this paper, under some restrictions on the initial data, we establish blow up of solutions with negative initial energy, using the same techniques as in [7].

Throughout this paper, $\|\cdot\|$ and $\|\cdot\|_p$ denote the usual $L^2(\Omega)$ norm and $L^p(\Omega)$ norm, respectively.

This paper is organized as follows. In Section 2, we state the local existence result. In Section 3, we show the blow up properties of solution in the case $p = q = 1$.

2. Local Existence

In this section, we state the local existence and uniqueness of the solution of problem (1.1).

Definition 1. A pair of functions (u, v) is said to be a weak solution of (1.1) on $[0, T]$ if $u, v \in C\left([0, T]; W_0^{1,2(m+1)}(\Omega) \cap L^{r+1}(\Omega)\right)$, $u_t \in C\left([0, T]; L^2(\Omega)\right) \cap L^{p+1}(\Omega \times (0, T))$ and $v_t \in C\left([0, T]; L^2(\Omega)\right) \cap L^{q+1}(\Omega \times (0, T))$. In addition, (u, v) satisfies

$$\begin{aligned} & \int_{\Omega} u'(t) \phi dx - \int_{\Omega} u_1(t) \phi dx + \int_{\Omega} \left(\rho(|\nabla u|^2) \nabla u\right) \nabla \phi dx \\ & \quad + \int_0^t \int_{\Omega} |u'|^{p-1} u' \phi dx d\tau \\ & = \int_0^t \int_{\Omega} f_1(u(\tau), v(\tau)) \phi dx d\tau, \end{aligned} \tag{2.1}$$

$$\begin{aligned} & \int_{\Omega} v'(t) \varphi dx - \int_{\Omega} v_1(t) \varphi dx + \int_{\Omega} \left(\rho(|\nabla v|^2) \nabla v\right) \nabla \varphi dx \\ & \quad + \int_0^t \int_{\Omega} |v'|^{q-1} v' \varphi dx d\tau \\ & = \int_0^t \int_{\Omega} f_2(u(\tau), v(\tau)) \varphi dx d\tau, \end{aligned} \tag{2.2}$$

for all test functions $\phi \in W_0^{1,2(m+1)}(\Omega) \cap L^{p+1}(\Omega)$, $\varphi \in W_0^{1,2(m+1)}(\Omega) \cap L^{q+1}(\Omega)$ and for almost all $t \in [0, T]$.

Now, we state the local existence theorem (see [3] for the proof).

Theorem 2. (*Local existence*). Assume (A1) holds. Then, for any initial data $u_0, v_0 \in W_0^{1,2(m+1)}(\Omega) \cap L^{r+1}(\Omega)$ and $u_1, v_1 \in L^2(\Omega)$, there exists a unique local weak solution (u, v) of problem (1.1) (in the sense of Definition 2) defined in $[0, T]$ for some $T > 0$, and satisfies the energy identity

$$E(t) + \int_0^t \left(\|u_{\tau}(\tau)\|_{p+1}^{p+1} + \|v_{\tau}(\tau)\|_{q+1}^{q+1} \right) d\tau = E(0), \tag{2.3}$$

where

$$E(t) = \frac{1}{2} \left(\|u_t\|^2 + \|v_t\|^2 \right) + \frac{1}{2} \int_{\Omega} \left(P(|\nabla u|^2) + P(|\nabla v|^2) \right) dx - \int_{\Omega} F(u, v) dx. \quad (2.4)$$

3. Blow Up of Solutions

In this section, we are going to consider the blow up of the solution for problem (1.1), when $p = q = 1$. For this purpose, we give the following lemma.

Lemma 2. (see [7]) *Suppose that $\psi(t)$ is a twice continuously differentiable function satisfying*

$$\begin{cases} \psi''(t) + \psi'(t) \geq C_0 \psi^{1+\alpha}(t), & t > 0, \\ \psi(0) > 0, & \psi'(0) \geq 0, \end{cases}$$

where $C_0 > 0$, $\alpha > 0$ are constants. Then, $\psi(t)$ blows up in finite time.

Theorem 3. *Let the assumptions of Theorem 3 hold. Assume further that $p = q = 1$. If initial data satisfies*

$$E(0) \leq 0, \quad \int_{\Omega} (u_0 u_1 + v_0 v_1) dx \geq 0,$$

then the corresponding solution blows up in finite time. In other words, there exists a positive constant T^* such that $\lim_{t \rightarrow T^*} (\|u\|^2 + \|v\|^2) = \infty$.

Proof. To apply Lemma 4, we define

$$\psi(t) = \frac{1}{2} \int_{\Omega} (|u|^2 + |v|^2) dx. \quad (3.1)$$

Therefore

$$\psi'(t) = \int_{\Omega} (uu_t + vv_t) dx, \quad (3.2)$$

and

$$\psi''(t) = \int_{\Omega} (u_t^2 + v_t^2) dx + \int_{\Omega} (uu_{tt} + vv_{tt}) dx. \quad (3.3)$$

Then, eq (1.1) is used to estimate (3.3) as follows

$$\begin{aligned}
 \psi''(t) &= \int_{\Omega} (u_t^2 + v_t^2) dx - \int_{\Omega} \left(\rho (|\nabla u|^2) |\nabla u|^2 + \rho (|\nabla v|^2) |\nabla v|^2 \right) dx \\
 &\quad - \int_{\Omega} (uu_t + vv_t) dx + (r+1) \int_{\Omega} F(u, v) dx \\
 &= \int_{\Omega} (u_t^2 + v_t^2) dx - b_1 \left(\|\nabla u\|^2 + \|\nabla v\|^2 \right) \\
 &\quad - b_2 \left(\|\nabla u\|_{2(m+1)}^{2(m+1)} + \|\nabla v\|_{2(m+1)}^{2(m+1)} \right) \\
 &\quad - \int_{\Omega} (uu_t + vv_t) dx + (r+1) \int_{\Omega} F(u, v) dx. \tag{3.4}
 \end{aligned}$$

Now, we exploit (3.4) to substitute for $\|\nabla u\|_{2(m+1)}^{2(m+1)} + \|\nabla v\|_{2(m+1)}^{2(m+1)}$; thus (3.4) takes the form

$$\begin{aligned}
 \psi''(t) + \psi'(t) &= (m+2) \left(\|u_t\|^2 + \|v_t\|^2 \right) + b_1 m \left(\|\nabla u\|^2 + \|\nabla v\|^2 \right) \\
 &\quad - 2(m+1) E(t) + (r-2m-1) \int_{\Omega} F(u, v) dx \\
 &\geq (m+2) \left(\|u_t\|^2 + \|v_t\|^2 \right) + b_1 m \left(\|\nabla u\|^2 + \|\nabla v\|^2 \right) \\
 &\quad - 2(m+1) E(t) + c_0 (r-2m-1) \left(\|u\|_{r+1}^{r+1} + \|v\|_{r+1}^{r+1} \right) \\
 &\geq \gamma \left(\|u\|_{r+1}^{r+1} + \|v\|_{r+1}^{r+1} \right), \tag{3.5}
 \end{aligned}$$

where $\gamma = c_0 (r - 2m - 1)$.

Now, Hölder inequality are used to estimates $\|u\|_{r+1}^{r+1}$ and $\|v\|_{r+1}^{r+1}$ as follows

$$\int_{\Omega} |u|^2 dx \leq \left(\int_{\Omega} |u|^{r+1} dx \right)^{\frac{2}{r+1}} \left(\int_{\Omega} 1 dx \right)^{\frac{r-1}{r+1}}.$$

W_n is called the volume of the domain Ω , then

$$\|u\|_{r+1}^{r+1} \geq \left(\int_{\Omega} |u|^2 dx \right)^{\frac{r+1}{2}} (W_n)^{-\left(\frac{r-1}{2}\right)}, \tag{3.6}$$

and similarly

$$\|v\|_{r+1}^{r+1} \geq \left(\int_{\Omega} |v|^2 dx \right)^{\frac{r+1}{2}} (W_n)^{-\left(\frac{r-1}{2}\right)}. \tag{3.7}$$

Consequently, we have

$$\psi''(t) + \psi'(t) \geq \gamma(W_n)^{-\left(\frac{r-1}{2}\right)} \left[\left(\int_{\Omega} |u|^2 dx \right)^{\frac{r+1}{2}} + \left(\int_{\Omega} |v|^2 dx \right)^{\frac{r+1}{2}} \right]. \quad (3.8)$$

In order to estimate the right-hand side in (3.8), we make use of the following inequality

$$(X + Y)^\rho \leq 2^{\rho-1} (X^\rho + Y^\rho),$$

$X, Y \geq 0, 1 \leq \rho < \infty$, applying the above inequality we have

$$2^{-\left(\frac{r-1}{2}\right)} \left(\int_{\Omega} |u|^2 dx + \int_{\Omega} |v|^2 dx \right)^{\frac{r+1}{2}} \leq \left(\int_{\Omega} |u|^2 dx \right)^{\frac{r+1}{2}} + \left(\int_{\Omega} |v|^2 dx \right)^{\frac{r+1}{2}}.$$

Consequently, (3.8) becomes

$$\begin{aligned} \psi''(t) + \psi'(t) &\geq 2^{-\left(\frac{r-1}{2}\right)} \gamma(W_n)^{-\left(\frac{r-1}{2}\right)} \left(\int_{\Omega} |u|^2 dx + \int_{\Omega} |v|^2 dx \right)^{\frac{r+1}{2}} \\ &= 2\gamma(W_n)^{-\left(\frac{r-1}{2}\right)} \psi^{\frac{r+1}{2}}(t). \end{aligned}$$

It is easy to verify that the requirements of Lemma 4 are satisfied by

$$C_0 = 2\gamma(W_n)^{-\left(\frac{r-1}{2}\right)} > 0 \text{ and } \alpha = \frac{r+1}{2} > 0.$$

Therefore $\psi(t)$ blows up in finite. □

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